The role of shocks in the InterStellar Medium

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The matter cycle in the galaxy



Quantitative view of the galactic cycle



Typical values

Velocities are ~ trans-sonic

=> Shocks are likely to form

	HIM	WNM	CNM	Diffuse	Dense	Discs	Sun
Density $\rho [\mathrm{cm}^{-3}]$	0.004	0.6	30	200	10^{4}	10^{10}	1 g.cm^{-3}
Temperature T [K]	3.10^{5}	5000	100	50	10	300	10^{6}
Length scale L [pc]	100	50	10	3	0.1	200 AU	$5.10^{-3} { m AU}$
Velocity $U [\rm km.s^{-1}]$	10	10	10	3	0.1	0.1	1
\mathcal{M}	0.2	2	13	7	0.5	0.1	0.02
$\mathcal{M}_\mathcal{G}$	130	20	15	6	0.8	0.08	0.003
\mathcal{R}^{-}	10^{2}	10^{5}	10^{7}	10^{7}	10^{6}	10^{9}	10^{17}
\mathcal{R}_m	10^{21}	10^{20}	10^{18}	10^{17}	10^{15}	10^{9}	10^{10}
\mathcal{R}_{AD}	10^{3}	10^{3}	10^{2}	10^{3}	10^{4}	10^{5}	10^{20}



Shocks through the galactic cycle



Shocks as probes of the cosmic cycle

Like a photographic developer, shocks uncover details of a medium otherwise cold, diffuse and sterile.



Outline

- Introduction on shocks physics The role of shocks for:
- Dissipative heat production
- Diffuse ISM chemistry
- Magnetic fields shaping
- I will skip: CR acceleration, dust processing





Shock wave



Shock wave



Shocks: bullet



Rankine Hugoniot

Flux conservation through a *steady* planar shock



Rankine Hugoniot

 Conservation of mass, momentum and magnetic flux in the steady shock *frame* induces relationships between preshock and post-shock physical conditions.

$$\begin{bmatrix} B_x \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho \ u_x \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} (B \times u)_y \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} (B \times u)_z \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho u_x u_y - B_x B_y \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho u_x u_z - B_x B_z \end{bmatrix}_{pre}^{post} = 0$$
$$\begin{bmatrix} \rho u_x^2 + P \end{bmatrix}_{pre}^{post} = 0$$

• Examples:

* Compression = Mach² in an isothermal shock * Max temperature $\sim u^2$ expresses conversion of kinetic to thermal energy in a viscous front $T_{\rm max} = 53 \,{\rm K} \,(u/1 \,{\rm km \ s^{-1}})^2$

For the molecular weight of the ISM:

Energy fluxes through a viscous ("J-type") shock in the ISM



The Paris-Durham shock code

Dissipative heating



Dissipation in the turbulent ISM Molecules, magnetic fields and Intermittency in coSmic Turbulence Following the energy trail...

Edith Falgarone

François Boulanger, Benjamin Godard, Pierre Hily-Blant, François Levrier, Pierre Lesaffre, Guillaume Pineau des Forêts,

Andrew Lehmann, Alba Vidal García, Thibaud Richard



2D turbulence: dissipation is very localised

Integrated Observables Line intensities in Polaris flare



Integrated Observables Centroid Velocity Increments



Shocks in 3D turbulence: dissipation is very localised



Pierre Lesaffre, G. Momferratos

Dissipation in decaying turbulence (incompressible runs)

 $n_{H} \sim 100/cm^{3}$ $<u^{2}>\sim<b^{2}/\rho>$ $Re=LU/v \sim 2.10^{7} 10^{3}$ $Re_{m}=LU/\eta \sim 2.10^{17} 10^{3}$ $Re_{AD}=L/U/t_{AD} \sim 10^{2}$

1 pc

Line of sight integrated dissipation:

$$\epsilon_{\rm diss} = \nu \rho S_{ij}[u] \partial_i u_j + \eta |\nabla \times B|^2 + F_{in}|u - v|^2$$

(Momferratos PhD thesis: 512³ spec. elts Incompressible simulations by ANK, pseudo-spectral code with AD)

Statistics of structures with strong dissipation





Momferratos et al. 2013

Dissipation in the diffuse ISM (compressible runs, Mach 4, isothermal)

 $n_{H} \sim 100/cm^{3}$ $<u^{2}>\sim<b^{2}/\rho>$ Re=LU/ $v \sim 2.10^{7} 10^{3}$ Re_m=LU/ $\eta \sim 2.10^{17} 10^{3}$ (1020³ pixels)

Heating nature in decaying MHD turbulence

~1 pc



Red: Ohmic, Green: Viscous shear, Blue: Viscous compression

(Momferratos PhD thesis: DUMSES simulations with careful treatment of viscous and resistive dissipation

Decaying turbulence (2D runs)



(way above average, but think intermittency)

Induced chemistry



Let's shake a small piece of ISM code CHEMSES=(RAMSES-AMR)+Paris Durham

Diffuse medium (nH=100, G_=1), viscous length resolved, 32 species followed

~ 0.003 pc



CPU time: 50 000 h at IDRIS (1 day on 2000 procs)

Simulated time: 10 000 yr.



Lesaffre et al. (in prep)

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H₂ excited and molecules produced by dissipation of 2D turbulence



G₀=1 Av=0.1

Find steady-state shocks (local fit of adiabatic fronts)



CO map produced in shocks (idem for OH, H₂O, CH, etc...)



Lesaffre et al. (in prep)

Fraction of some hydrides in background and shocked region

Both regions have same volume, but shocked regions have more molecules



1D steady models can retrieve the molecular yields (CPU time=1h)



Prospects

Intermittent statistics of the dissipation



3D simulations (cf Momferratos et al. 2013)

Dissipation strength => Molecules Formation + excitation

Molecular yields from Shocks (for example)



1D simulations

simulations

CO map

Magnetic fields processing



B orientation in oblique shocks



Figure 2. The effect on magnetic field orientation of the three classes of MHD shock waves. Fast shocks (left) increase the angle between the field and shock normal, intermediate shocks (middle) reverse the sign of the angle and slow shocks (right) decrease it. Hence the magnetic field strengthens across fast shocks and weakens across slow shocks.

In the 'SHOCKFIND' paper by Lehmann, Federrath, Wardle (201

LOFAR & Planck data

Rotation measure (Rm= Σ ne.Bz.dz) overlayed with B field direction



Dissipative structures extraction

Find connected sets where dissipation > mean + 2.standard deviations



B field is mostly parallel to structures

Thibaud Richard



Conclusions

- Shocks are everywhere
- They convert kinetic energy into thermal energy
- They are good probes of the dynamics and chemistry of the gas
- They allow to make good use of up to date observational capabilities
- Still a lot of fun ahead !



Current Developments of the Paris-Durham shock code

- 3D models (Le Ngoc Tram) Tram+2018
- Stellar winds (Le Ngoc Tram) Tram+2018 in prep.
- Line post-processing (A. Gusdorf + Tram)
- Irradiated / self-irradiated shocks (A. Lehmann, B. Godard) Godard+2018 in prep.
- Oblique shocks (A. Lehmann, P.L.)

Stay Tuned !



Thanks !



Integrated Observables

<u>Centroid velocity:</u> first moment of the l.o.s. velocity

$$\mathrm{CV}(x,y) = \int_0^L u_z(x,y,z) \ dz$$

Assuming that total dissipation powers the line (or that a chemical tracer appears right where there is heating):

$$\operatorname{CV}_w(x,y) = \frac{1}{\langle \varepsilon \rangle} \int_0^L \varepsilon(x,y,z) \ u_z(x,y,z) \ dz$$

Other variables: Stokes parameters of the polarization (Q, U, I, P)... (assuming grains are perfectly aligned to local B)

Observable increments vs. dissipation Lbox / 64



• <u>Background</u>: Dissipation rates Ohmic Viscous AD

• <u>Contours:</u> Increments of

integrated observables:

- LOS velocity (white)
- Stokes Q (green)
- Stokes U (red)
- POS polarisation angle (blue)

<u>NOTE:</u> different observables trace different parts of the dissipative structures

Observable increments vs. dissipation Lbox / 8



• <u>Background</u>: Dissipation rates Ohmic Viscous AD

• <u>Contours:</u> Increments of

integrated observables:

- LOS velocity (white)
- Stokes Q (green)
- Stokes U (red)
- POS polarisation angle (blue)

<u>NOTE:</u> increment of polarisation angle (blue contours) are less correlated to dissipation. Better use Q,U.

Observable increments vs. dissipation Lbox/2



• <u>Background</u>: Dissipation rates Ohmic Viscous AD

• <u>Contours:</u> Increments of

integrated observables:

- LOS velocity (white)
- Stokes Q (green)
- Stokes U (red)
- POS polarisation angle (blue)

<u>NOTE:</u> increment of polarisation angle (blue contours) are less correlated to dissipation. Better use Q,U.